



Born Approximation

M.Sc. 4th Semester

MPHYEC-1: Advanced Quantum Mechanics

Unit I

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Lecture Outline

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- Green's functions
- Scattering amplitude
- General solution of Schrödinger eq. in terms of Green's function
- Born Series
- Assignment

Green's functions

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The Green's function is obtained by solving the point source equation:

$$(\nabla^2 + k^2)G(\vec{r} - \vec{r}') = \delta(\vec{r} - \vec{r}') \quad \text{.....1}$$

where $G(\vec{r} - \vec{r}')$ and $\delta(\vec{r} - \vec{r}')$ are given by their Fourier transforms as follows:

$$G(\vec{r} - \vec{r}') = \frac{1}{(2\pi)^3} \int e^{i\vec{q}\cdot(\vec{r}-\vec{r}')} \tilde{G}(\vec{q}) d^3q, \quad \text{.....2}$$
$$\delta(\vec{r} - \vec{r}') = \frac{1}{(2\pi)^3} \int e^{i\vec{q}\cdot(\vec{r}-\vec{r}')} d^3q.$$

Substituting Eqn (2) into Eqn(1) gives

$$(-\vec{q}^2 + k^2) \tilde{G}(\vec{q}) = 1 \quad \implies \quad \tilde{G}(\vec{q}) = \frac{1}{k^2 - \vec{q}^2}. \quad \text{.....3}$$

The expression for $G(\vec{r} - \vec{r}')$ can be obtained by inserting (3) into (2)

$$G(\vec{r} - \vec{r}') = \frac{1}{(2\pi)^3} \int \frac{e^{i\vec{q}\cdot(\vec{r}-\vec{r}')}}{k^2 - q^2} d^3q. \quad \dots\dots 4$$

$$G(\vec{r} - \vec{r}') = \frac{1}{(2\pi)^3} \int_0^\infty \frac{q^2 dq}{k^2 - q^2} \int_0^\pi e^{iq|\vec{r}-\vec{r}'|\cos\theta} \sin\theta d\theta \int_0^{2\pi} d\varphi, \quad \dots\dots 5$$

To integrate over angle in (5) we need to make the variable change $x = \cos\theta$

$$\int_0^\pi e^{iq|\vec{r}-\vec{r}'|\cos\theta} \sin\theta d\theta = \int_{-1}^1 e^{iq|\vec{r}-\vec{r}'|x} dx = \frac{1}{iq|\vec{r}-\vec{r}'|} \left(e^{iq|\vec{r}-\vec{r}'|} - e^{-iq|\vec{r}-\vec{r}'|} \right) \quad \dots\dots 6$$

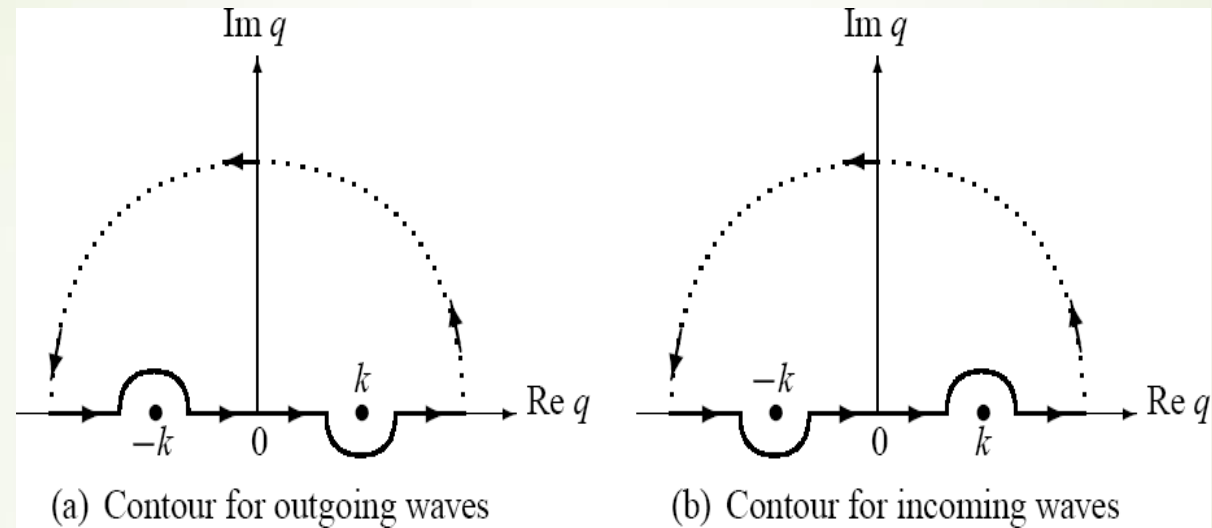
Thus, (4) becomes

$$G(\vec{r} - \vec{r}') = \frac{1}{4\pi^2 i |\vec{r} - \vec{r}'|} \int_0^\infty \frac{q}{k^2 - q^2} \left(e^{iq|\vec{r}-\vec{r}'|} - e^{-iq|\vec{r}-\vec{r}'|} \right) dq, \quad \dots\dots 7$$

or

$$G(\vec{r} - \vec{r}') = -\frac{1}{4\pi^2 i |\vec{r} - \vec{r}'|} \int_{-\infty}^{+\infty} \frac{q e^{iq|\vec{r}-\vec{r}'|}}{q^2 - k^2} dq.$$

The integral in (7) can be evaluated by the method of residues by closing the contour in the upper half of the q -plane:



The integral is equal to $2\pi i$ times the residue of the integrand at the poles.

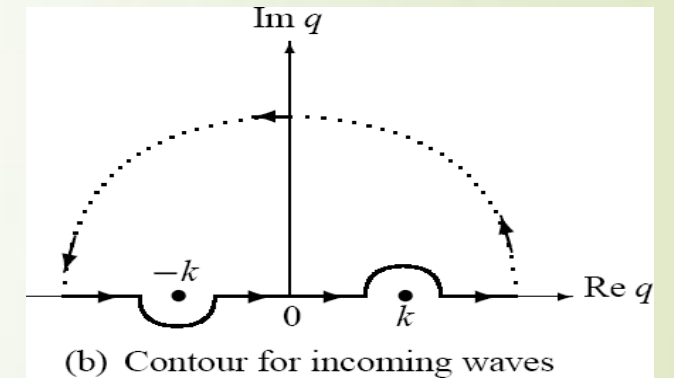
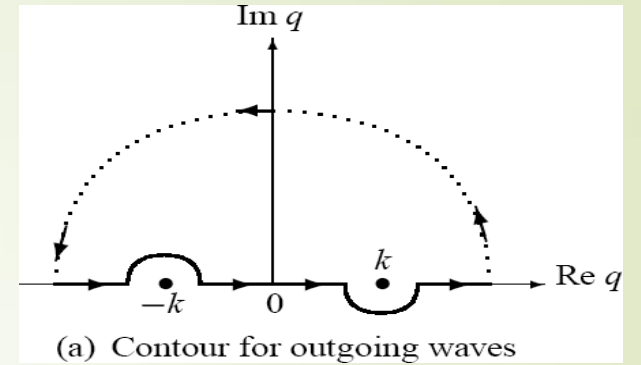
Since there are two poles, $q = \pm k$, the integral has two possible values:

- the value corresponding to the pole at $q = k$, which lies inside the contour of integration in Figure 1a, is given by

$$G_+(\vec{r} - \vec{r}') = -\frac{1}{4\pi} \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \quad \dots\dots 8$$

- the value corresponding to the pole at $q = -k$, Figure 1b, is

$$G_-(\vec{r} - \vec{r}') = -\frac{1}{4\pi} \frac{e^{-ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \quad \dots\dots 9$$



Green's function $G_+(\vec{r} - \vec{r}')$ represents an *outgoing spherical wave* emitted from r' and the function $G_-(\vec{r} - \vec{r}')$ corresponds to an *incoming wave* that converges onto r .

Since the scattered waves are outgoing waves, only $G_+(\vec{r} - \vec{r}')$ is of interest to us.

Scattering amplitude

We are going to show here that we can obtain the **differential cross section in the CM frame** from an **asymptotic form** of the solution of the **Schrödinger equation**:

$$-\frac{\hbar^2}{2\mu} \nabla^2 \psi(\vec{r}) + \hat{V}(r) \psi(\vec{r}) = E \psi(\vec{r}) \quad \text{.....10}$$

□ Let us first **focus on the determination of the scattering amplitude $f(\theta, \varphi)$** , it can be obtained from the solutions of (10), which in turn can be rewritten as

$$(\nabla^2 + k^2) \psi(\vec{r}) = \frac{2\mu}{\hbar^2} V(\vec{r}) \psi(\vec{r}), \quad \text{where } k^2 = \frac{2\mu E}{\hbar^2} \quad \text{.....11}$$

The general solution of the equation (11) consists of a sum of **two components**:

1) a general solution to the **homogeneous equation**:

$$(\nabla^2 + k_0^2) \psi_{\text{homo}}(\vec{r}) = 0, \quad \text{where } k_0^2 = 2\mu E / \hbar^2. \quad \text{.....12}$$

In (12) $\psi_{\text{homo}}(\vec{r}) \rightarrow \phi_{\text{inc}}(\vec{r}) = A e^{i\vec{k}_0 \cdot \vec{r}}$, is the **incident plane wave**

2) and a particular solution of (11) with the **interaction potential**

General solution of Schrödinger eq. in terms of Green's function

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The **general solution** of (11) can be expressed in terms of **Green's function**.

$$\psi(\vec{r}) = \phi_{inc}(\vec{r}) + \frac{2\mu}{\hbar^2} \int G(\vec{r} - \vec{r}') V(\vec{r}') \psi(\vec{r}') d^3r', \quad \dots\dots 13$$

where $\phi_{inc}(\vec{r}) = e^{i\vec{k}_0 \cdot \vec{r}}$ and $G(\vec{r} - \vec{r}')$ is the **Green's function** corresponding to the operator on the left side of eq.(12)

Inserting (8) into (13) we obtain for the total scattered wave function:

This is an **integral equation**.

All we have done is to rewrite the Schrödinger (differential) equation (10) in an **integral form** (13), which is more suitable for scattering theory.

Note that (13) can be **solved approximately** by means of a series of successive or **iterative** approximations, known as the **Born series**.

Born Series

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- the **zero-order solution** is given by $\psi_0(\vec{r}) = \phi_{inc}(\vec{r})$
- the **first-order solution** $\psi_1(\vec{r})$ is obtained by inserting $\psi_0(\vec{r}) = \phi_{inc}(\vec{r})$ into the integral of (13):

$$\begin{aligned}\psi_1(\vec{r}) &= \phi_{inc}(\vec{r}) - \frac{\mu}{2\pi\hbar^2} \int \frac{e^{ik|\vec{r}-\vec{r}_1|}}{|\vec{r}-\vec{r}_1|} V(\vec{r}_1) \psi_0(\vec{r}_1) d^3r_1 \\ &= \phi_{inc}(\vec{r}) - \frac{\mu}{2\pi\hbar^2} \int \frac{e^{ik|\vec{r}-\vec{r}_1|}}{|\vec{r}-\vec{r}_1|} V(\vec{r}_1) \phi_{inc}(\vec{r}_1) d^3r_1.\end{aligned}\tag{14}$$

- the **second order solution** is obtained by inserting $\psi_1(\vec{r})$ into (13):

$$\begin{aligned}\psi_2(\vec{r}) &= \phi_{inc}(\vec{r}) - \frac{\mu}{2\pi\hbar^2} \int \frac{e^{ik|\vec{r}-\vec{r}_2|}}{|\vec{r}-\vec{r}_2|} V(\vec{r}_2) \psi_1(\vec{r}_2) d^3r_2 \\ &= \phi_{inc}(\vec{r}) - \frac{\mu}{2\pi\hbar^2} \int \frac{e^{ik|\vec{r}-\vec{r}_2|}}{|\vec{r}-\vec{r}_2|} V(\vec{r}_2) \phi_{inc}(\vec{r}_2) d^3r_2 \\ &\quad + \left(\frac{\mu}{2\pi\hbar^2}\right)^2 \int \frac{e^{ik|\vec{r}-\vec{r}_2|}}{|\vec{r}-\vec{r}_2|} V(\vec{r}_2) d^3r_2 \int \frac{e^{ik|\vec{r}_2-\vec{r}_1|}}{|\vec{r}_2-\vec{r}_1|} V(\vec{r}_1) \phi_{inc}(\vec{r}_1) d^3r_1.\end{aligned}\tag{15}$$

continuing in this way, we can obtain $\psi(\vec{r})$ to any desired order

Assignment

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1. Describe the method of Green's function to solve scattering Hamiltonian.
2. Explain the Green's function for point source.
3. Explain Born series expansion.



Thank You